

SOFIA FIFI-LS Observations of Sgr B1: Ionization Structure and Sources of Excitation

Janet P. Simpson¹⁽¹⁾, Sean W. J. Colgan², Angela S. Cotera¹, Michael J. Kaufman³⁽¹⁾, and Susan R. Stolovy⁴

SETI Institute, 189 Bernardo Avenue, Mountain View, CA 94043, USA; janet.p.simpson@gmail.com

NASA Ames Research Center, MS 245-6, Moffett Field, CA 94035-1000, USA

³ San Jose State University Department of Physics and Astronomy, One Washington Square, San Jose, CA 95192-0106, USA

⁴ El Camino College Physics Department, 16007 Crenshaw Boulevard, Torrance, CA 90506, USA

Received 2018 August 11; revised 2018 October 4; accepted 2018 October 15; published 2018 October 30

Abstract

The current paradigm of Galactic Center (GC) gas motions and star formation envisions sequential star formation in streams of gas as they pass near the supermassive black hole Sgr A^{*}. This is based on the relative positions of dense molecular clouds, the very young star-forming region Sgr B2, the much older region Sgr C, and the several Myr old Arches and Quintuplet Clusters. Because Sgr B1 is found with Sgr B2 in a common envelope of molecular gas and far-infrared emission, the two sources are thought to be physically related, even though there are indicators of a significantly greater age for Sgr B1. To clarify the status of Sgr B1, we have mapped it with the FIFI-LS spectrometer on the Stratospheric Observatory for Infrared Astronomy in the far-infrared lines of [O III] 52 and 88 μ m. From the ratios of these lines and lines measured with the *Spitzer* Infrared Spectrograph, we find that there are at least eight separate sub-regions that must contain the stars that excite the gas. We infer spectral energy distributions (SEDs) of the ionizing sources from models and find they are in agreement only with SEDs of late O stars augmented at the highest frequencies with interstellar X-rays from fast shocks. We suggest that although the gas, from its velocity structure, must be part of the very young Sgr B2 complex, the stars that are ionizing the gas were not formed there but are the remnants of a previous generation of star formation in the GC.

Key words: Galaxy: center - H II regions - infrared: ISM - ISM: individual objects (Sgr B1) - X-rays: ISM

1. Introduction

A recent model of star formation in the Galactic Center (GC) posits that there are streams of gas in open orbits around the nucleus (azimuthal period 3.69 Myr), with stars forming when the gas is compressed as it passes the pericenter, Sgr A (Longmore et al. 2013; Kruijssen et al. 2015). In particular, for the dense molecular clouds at positive Galactic longitudes that can be described as lying on the front side of the orbit, the amount of star formation increases as a function of distance from Sgr A (and hence time in the orbital model), with the most active star formation occurring in the very young Sgr B2 region (Longmore et al. 2013). Farther along the orbit at negative Galactic longitude is found another massive but much older GC H II region, Sgr C. The Arches and Quintuplet Clusters formed from gas clouds that passed pericenter at even earlier times (Kruijssen et al. 2015).

An anomaly in this scenario, however, is the luminous H II region Sgr B1, which appears to be part of the same giant molecular cloud as Sgr B2 but is already showing signs of dispersal as though due to winds from much earlier star formation (e.g., Mehringer et al. 1992, hereafter M92). In the Kruijssen et al. (2015) model, Barnes et al. (2017) found that Sgr B1 is on the back side of the orbit at an age of 1.5 Myr, versus 0.7 Myr for Sgr B2.

Even this age may be too young for Sgr B1. Recently, Simpson (2018, hereafter S18) estimated an age of 4.6 Myr for Sgr B1 from her inferred shape of the spectral energy distribution (SED) of the stars that ionize this H II region (SEDs computed with Starburst99; Leitherer et al. 2014). In S18, all the mid-infrared spectra of the GC taken by the Infrared Spectrograph with *Spitzer Space Telescope* were reanalyzed, and the observed line ratios were compared to H II region models computed with Cloudy 17.00 (Ferland et al. 2017). The results

were that the Quintuplet Cluster region, the Arched Filaments, Sgr B1, and Sgr C H II regions generally have the low excitation predicted by models ionized by SEDs of ages ranging from 2.5 to 5 Myr. The data emphasized in the comparison were the [S III] 33/[Si II] 34 μ m line ratios, the [O IV] 26/[S III] 33 μ m line ratios, and the [Ne III] 15.6/[S III] 18.7 μ m line ratios. The first, a function of the ionization parameter (photon density divided by the electron density) is an indicator of the dilution of the radiation field, and hence shows the relative closeness of the ionizing stars. The second ratio, because O^{3+} has an ionization potential (IP) of 54.9 eV, shows the high-energy content of the SEDs; S18 concluded that the SEDs ionizing GC HII regions have X-rays additional to the Starburst99 high-energy photons. For the last, because the shape of the SED > 13.6 eV is a strong function of the stellar effective temperatures (T_{eff}) (or age, for Starburst99 models), ratios of the Ne^{++} and S^{++} ions, with IP equal 41 and 23 eV, respectively, indicate this shape and hence can be used to infer $T_{\rm eff}$ or age.

In this Letter, we describe observations of Sgr B1 produced in lines of O⁺⁺ (which has an intermediate IP of 35 eV), made with the Field Imaging Far-Infrared Line Spectrometer (FIFI-LS; Colditz et al. 2018; Fischer et al. 2018) on the Stratospheric Observatory for Infrared Astronomy (SOFIA; Young et al. 2012; Temi et al. 2014). Section 2 describes the observations, Section 3 depicts the computed electron densities (N_e), and the ratios of the ions O⁺⁺/S⁺⁺ and Ne⁺⁺/O⁺⁺, where the [Ne III] and [S III] line intensities were taken from S18. Section 4 compares the ionic abundance ratios and line ratios to models of H II regions and shocks, and Section 5 presents the summary and conclusions. We will conclude that Sgr B1 is ionized by widely dispersed and relative cool OB stars with SEDs including X-rays, possibly from high-velocity shocks. THE ASTROPHYSICAL JOURNAL LETTERS, 867:L13 (6pp), 2018 November 1



Figure 1. Observed [O III] line intensities in Sgr B1. All but the darkest blue pixels have S/N > 3. The black contours are the 8.4 GHz Very Large Array (VLA) intensities from M92 and the cyan contours are the 70 μ m *Herschel* Hi-GAL intensities (Molinari et al. 2011, 2016). (a) The [O III] 52 μ m line (maximum intensity = 1.92×10^{-16} W m⁻² arcsec⁻²). (b) The [O III] 88 μ m line (maximum intensity = 1.57×10^{-16} W m⁻² arcsec⁻²).

2. Observations

Sgr B1 was mapped with FIFI-LS in 2016 and 2017 July, with SOFIA flying from Christchurch, NZ. FIFI-LS has two spectrometers that operate simultaneously. The [O III] lines at 51.81 and 88.36 μ m were observed with the blue channel (6" pixels, spectral resolution ~930 and ~600, respectively). The red channel lines ([O I] 145.53 μ m and [C II] 157.74 μ m) will be discussed in a later paper.

The parts of the map with the brightest radio emission were observed in 2016 and the rest in 2017. In 2016, the chopper throw was 4' at position angles of $133^{\circ}-155^{\circ}$ east from north, approximately perpendicular to the Galactic plane. In 2017, the chopper position angle for all observations was 135° with chopper throws 5' or 6' for the 52 μ m line and 6!7 for the 88 μ m line. All lines were measured in the "nod match chop" mode (telescope boresight halfway between the source and the reference positions); the resulting beam pattern is the diffraction-limited beam size (wavelength divided by 12; Young et al. 2012) plus about 1" of coma added in quadrature for each arcmin of chopper throw.

The FIFI-LS detectors form a grid of 5×5 pixels on the sky; sources are dithered with 3" steps and then the whole array is stepped by about 30" for another mini-map. The FIFI-LS pipeline combines all of the observations taken consecutively into a cube (R.A., decl., and wavelength) with 1" pixels. To combine the resulting nine sub-cubes per line into a single cube, a large cube was first defined that covered the entire region as seen in the red channel in R.A. and decl. The wavelength dimension for this cube was that of the sub-cube with the maximum number of wavelength values (each of the sub-cubes had a slightly different wavelength scale). Then the spectra of each of the sub-cubes were interpolated onto this wavelength array (because the FIFI-LS line profiles are greatly oversampled, the interpolated fluxes were usually quite similar to the fluxes of their neighboring points). Next, all of the sub-cubes were shifted in R.A. and decl. onto the big grid. Because no interpolation was used, such shifts could be in error by as much as half a pixel (0.75), which is small compared to the $7^{\prime\prime}-10^{\prime\prime}$ spatial resolution. Finally, all of the new large cubes were combined—where there were multiple integrations on a single point on the sky, the spectra were averaged and the errors were combined in quadrature.

Line intensities were estimated for each spectrum, pixel by pixel on the sky, by integrating the line profile over a given fixed range with a few pixels on each end defining the continuum. Typically, the number of wavelength pixels for the line was much larger than the number for the continuum, especially for the 52 μ m line, which has a deep telluric H₂O line very close on the short wavelength side. The uncertainties for the line fluxes were estimated from the rms deviation of the data from a fitted line-profile function, which gave a good representation of the overall shape of the spectrum but an unreliable flux. Because there are thousands of pixels, this fitting had to be done with an automated line-fitting program with little hand checking, except for those measurements with signal-to-noise ratio (S/N) < 4, which were checked by eye; very noisy line measurements were then rejected.

Maps of the 52 and 88 μ m line intensities (linear scale) are plotted in Figure 1. The measured continuum values are not useful because the small chopper throw in 2016 resulted in the telescope chopping onto and subtracting extra-source continuum (although not line, as estimated from the *Spitzer* maps of S18). In this figure, we see that the [O III] morphology is distinctly different from both the 8.4 GHz radio map (proportional to the emission measure, $\int N_e^2 dl$) and the 70 μ m Herschel image (proportional to the column density of warm dust), which are quite similar.



Figure 2. Results from ionic abundance ratio calculations. The contours in each map are those of the 8.4 GHz VLA image from M92. (a) Computed electron densities from the [O III] 52/[O III] 88 μ m line ratios. The colored areas are where both [O III] lines were observed and the line ratios have S/N \ge 4. (b) The [S III] 33/[Si II] 34 μ m line ratio map from S18 (demonstrating the incomplete *Spitzer* map coverage). The locations of the eight sub-regions discussed in Section 4.1 are plotted as the numbered black boxes. (c) The ionic O⁺⁺/S⁺⁺ ratio. (d) The ionic Ne⁺⁺/O⁺⁺ ratio. The large black stars show the locations of the O supergiant and the Wolf–Rayet (WR) stars identified by Mauerhan et al. (2010).

3. Results

We corrected the observed line intensities for the extinction estimated by S18 and computed ionic abundance ratios using the method described by Simpson et al. 2004, 2007, 2012), with O⁺⁺ collision strengths from Storey & Sochi (2015) and other atomic physics from Table 1 of S18. For estimates of N_e , we required that the 52/88 μ m line ratio was measured with S/N \ge 4 and required line measurements with S/N \ge 3 for ionic abundance ratios. Maps of these results (all linear scale) are plotted in Figure 2.

Figure 2(a) shows a plot of N_e with the radio contours overplotted. The average density for the high S/N region is low, $N_e \sim 300 \text{ cm}^{-3}$. The regions with the highest intensity radio contours on the western side of Sgr B1—the "Ionized Bar" $(17^{h}47^{m}01^{s}4-28^{\circ}30'37'')$ and the "Ionized Rim" $(17^{h}46^{m}59^{s}9-28^{\circ}32'53'')$ of M92—are examples of low densities. Here, a substantial distance of ionized gas along the line of sight is required to produced the observed emission measure, given the low density. Low densities for this region are also estimated from the ratios of the lower-excitation [S III] 19/[S III] 33 μ m lines, where S⁺⁺ is the dominant ionization state of sulfur (S18). We speculate that these structures are edge-on sheets of gas surrounding local holes or bubbles blown by winds from hot stars. The eastern parts of Sgr B1, however, have densities commensurate with the intensity of the radio emission. The highest densities are found in regions of low intensity and may be external to Sgr B1.

Figure 2(b) shows the [S III] $33/[Si II] 34 \mu m$ line ratios from *Spitzer* (S18). This ratio is a function of the photon density of

 Table 1

 Positions Inferred to have nearby Sources of Ionization

Position	R.A. (J2000)	Decl. (J2000)	$\frac{\log N_{\rm Lyc}{}^{\rm a}}{({\rm s}^{-1})}$	$\binom{N_e}{(\mathrm{cm}^{-3})}$	$0^{++}/S^{++}$	Ne ⁺⁺ /O ⁺⁺	$\frac{[O III]88 \mu\text{m}}{[S III]33 \mu\text{m}}$ Line Ratio	[Ne III] 15.6 μm [Ο III] 88 μm Line Ratio
1	17 46 52.00	-28 32 41	48.09	841	6.4	0.50	0.16	0.89
2	17 46 56.05	-28 29 51	48.27		4.1	0.24	0.31	0.28
3	17 46 58.40	-28 29 55	48.29	269	5.7	0.21	0.38	0.30
4	17 47 00.15	-28 30.25	48.72	723	4.4	0.08	0.06	0.18
5	17 47 02.65	-28 30 48	48.70	297	7.5	0.09	0.17	0.20
6	17 47 04.00	-28 31 18	48.58	181	4.9	0.18	0.30	0.26
7	17 47 06.90	$-28 \ 31 \ 03$	48.73	342	8.7	0.19	0.34	0.42
8	17 47 11.98	-28 31 25	48.62	352	6.6	0.17	0.38	0.32

Note.

^a The estimated numbers of photons required to ionize the sub-regions were derived using Equation (4) of Simpson et al. (2012), where the measured [S III] 33 μ m line flux was integrated over the sub-region, the assumed S/H ratio was 1.90×10^{-5} , and the assumed S⁺⁺/S ratio was 0.8 (a likely value for these low-excitation H II regions, which have much lower and more uncertain O⁺⁺/O). The total estimated N_{Lyc} from [S III] is 3.3×10^{50} s⁻¹, in agreement with the estimated number of ionizing photons from the radio, 3×10^{50} s⁻¹ (S18), for assumed distance 8 kpc.

the radiation field, and hence demonstrates that the exciting stars are widespread, with occasional large gaps (ratio ≤ 1.5) where there are few ionizing stars.

Figures 2(c) and (d) show the O^{++}/S^{++} and Ne^{++}/O^{++} ratios, respectively (the O^{++}/S^{++} ratio has better coverage than the Ne^{++}/O^{++} ratio because we could use both the [S III] 19 and 33 μ m lines). These ratios are indicative of the shapes of the ionizing SEDs and so are an indication of the locations of the exciting stars with the highest T_{eff} .

4. Discussion

4.1. Sources of Ionization

We immediately notice that Sgr B1 is not ionized by a central star cluster—there are multiple regions of higher ionization, as seen in the O^{++}/S^{++} and Ne^{++}/O^{++} ratios plotted in Figure 2(c) and (d). After smoothing the O^{++}/S^{++} ratio by 30", we selected the eight regions with highest O^{++}/S^{++} ; they are delineated in Figure 2(b) and listed in Table 1. There are undoubtedly additional sub-regions, but the combined *Spitzer* and SOFIA data do not have enough coverage (e.g., Figure 2) to adequately define them.

It is particularly interesting that these regions of higher ionization do not coincide with the peaks in the radio emission (contours in Figures 1 and 2). M92 measured the flux density in each of these peaks and estimated the ZAMS spectral types of the OB stars that would be required to ionize them. Our observations, which mostly locate the ionizing stars in volumes separate from the ionized gas, show that the ionizing stars must be positioned in some pattern other than at the peaks of the radio emission. For example, the "Ionized Rim" requires the most ionizing photons; it certainly does not have the excitation corresponding to M92's suggested O6 star. Multiple stars of cooler temperature would be needed to ionize this region.

The morphology of the extended features and shell structures in the radio emission led M92 to suggest that Sgr B1 is an evolved H II region. There certainly are no dense, ultracompact H II regions like there are in Sgr B2 (e.g., Mehringer et al. 1993; De Pree et al. 2015), which is probably physically associated owing to the similar velocities and apparent gas structures that bridge the two regions (M92).

If Sgr B1 is indeed an evolved H II region, it should contain stars that have evolved significantly. Two to three Wolf–Rayet (WR) stars and an O supergiant were found in Sgr B1 by Mauerhan et al. (2010) and are plotted in Figure 2(d). The presence of evolved stars in the line of sight to Sgr B1 might not be significant—Habibi et al. (2014) simulated orbits for stars drifting away from the Quintuplet and Arches Clusters (ages 4.8 ± 1.1 and 3.5 ± 0.7 Myr, Schneider et al. 2014). They found that some of their simulated stars could travel as far from the Quintuplet Cluster as Sgr B1.

Another possibility could be that while the ionized and molecular gas is part of the Sgr B molecular cloud (M92; Mehringer et al. 1995; Lang et al. 2010), the ionizing stars originate in a much older cluster that has already orbited once around Sgr A, much as is thought to have occurred for the Quintuplet and Arches Clusters (e.g., Habibi et al. 2014; Kruijssen et al. 2015). These stars light up the edges of local molecular clouds (Lang et al. 2002) but are not near the locations of their formation (Stolte et al. 2014). We note that the multiple young stellar objects found in Sgr B1 (An et al. 2011, 2017), none especially luminous, are all found in the denser of the radio peaks; we suggest that these may be indicators of star formation triggered by the outflows from the very luminous ionizing stars.

4.2. Ionizing SEDs

To learn more about the ionizing SEDs in Sgr B1, we compared the line ratios and ionic ratios (ionic ratios being more reliable because they account for variations in density) observed in the selected sub-regions (Table 1 and Figure 2) to grids of H II region models. By plotting Ne⁺⁺/O⁺⁺ versus O⁺⁺/S⁺⁺ (ion or line ratio), we sample the 23–63 eV portion of the extreme ultraviolet SED as found in both the stars that ionize H II regions and the stellar atmosphere models used in H II region models. Such comparisons inform us both about the $T_{\rm eff}$ of the H II region's exciting stars and about the completeness of the codes used to compute the stellar atmosphere model SEDs (Simpson et al. 2004; Rubin et al. 2008). The averages of the sub-regions (Table 1) are plotted in Figure 3, along with the ratios from the models.

These models include the coarse grid of H II region models from S18. Her models all had hydrogen densities $N_p = 100 \text{ cm}^{-3}$, ionizing photon luminosities of 10^{50} s^{-1} , GC abundances (S18), filling factors ranging from 0.001 to 1.0, H II region inner radii ranging from 1 to 10 pc, and ionizing stellar SEDs from Starburst99 (Leitherer et al. 1999, 2014) with ages



Figure 3. Ionic and line ratios compared to models. In both panels, the black dots are our observed positions as described in the text and listed in Table 1. The asterisks are models from S18 computed with Cloudy (Ferland et al. 2017) with O/ S = 36 and Ne/O = 0.25, ionizing SEDs from Starburst99 (Leitherer et al. 2014), and extra X-rays represented by blackbodies with $T_{\rm BB} = 10^{6.5}$ K, where the colors signify the ages of the Starburst99 model SEDs: blue is $10^{6.7}$ years, green is $10^{6.65}$ years, yellow is $10^{6.6}$ years, yellow is $10^{6.6}$ years, red-orange is $10^{6.4}$ years, red is $10^{6.3}$ years, magenta is $10^{6.2}$ years, and purple is $10^{6.0}$ years. Squares are newly computed models with input density $= 300 \text{ cm}^{-3}$, the same Starburst99 SEDs and otherwise the same input parameters but no X-rays, and the same color scheme. Upward- and downward-pointed triangles are similar models with supergiant and dwarf SEDs computed with WM-BASIC (Pauldrach et al. 2001) taken from the Binary Population and Spectral Synthesis (BPASS) project (Eldridge et al. 2017); the colors represent $T_{\rm eff}$ equal to blue: 30,200 K, green: 32,300 K, orange: 34,600 K, red-orange: 37,200 K, and red: 40,000 K. For all models, the larger and smaller symbols represent models whose predicted [S III] 33/[Si II] 34 μm line ratios are within or outside the range of the observed ratio, 1.5–3.0, respectively. (a) The Ne⁺⁺/O⁺⁺ ratio vs. the O⁺⁺/S⁺⁺ ratio. (b) The [Ne III] 15.6/[O III] 88 μ m line ratio vs. the [O III] 88/[S III] 33 μ m line ratio, where the line intensities were corrected for extinction from S18. Recomputing the models plotted with asterisks for a higher $N_p = 300 \text{ cm}^{-3}$ would have them move in the direction of the blue arrow. Following Ho et al. (2014), the cyan and brown lines are the result of adding shock emission from Allen et al. (2008; see the text). The gray crosses are combinations of models plus shocks that predict the [S III] 33/[Si II] 34 μm line ratio within the observed range.

ranging from 10^6 to $10^{6.7}$ years and augmented with X-rays. For computational ease, the X-rays were represented by blackbodies with temperatures, $T_{\rm BB}$, of either $10^{6.0}$ or $10^{6.5}$ K and blackbody luminosities, $L_{\rm BB}$, ranging from 10^{37} to 10^{39} erg s⁻¹. However, because the line that required the X-rays for fitting was the [O IV] 26 μ m line with IP 54.9– 77 eV, only the X-rays in this energy range were significant to the model fits. The integrated ionizing-luminosity from 55 to 77 eV is 1.9×10^{35} or 6.2×10^{35} erg s⁻¹ for the two bestfitting models for Sgr B1 (input 3×10^{50} s⁻¹ ionizing photons, $T_{\rm BB} = 10^{6.5}$ K, and $L_{\rm BB} = 3 \times 10^{38.0}$ or $3 \times 10^{38.5}$ erg s⁻¹, respectively). Ratios from the models augmented with X-rays are plotted with asterisks in Figure 3 (omitting all models with predicted [S III] 33/[Si II] 34 μ m line ratios < 0.75).

Additional models with no X-rays were computed for this paper using Cloudy 17.01 (Ferland et al. 2017). The models using stellar SEDs from Starburst99 have the same variation in SED age, inner radius, and filling factor as the models with X-rays but $N_p = 300 \text{ cm}^{-3}$. These models are plotted with squares in Figure 3, where it is seen that none of these models has line ratios that lie in the region of the observed line ratios. We conclude that the only models that use Starburst99 SEDs that also have X-rays agree with the data.

However, the SEDs from Starburst99 are the summation of the individual stellar SEDs from a massive cluster of stars, starting with an initial mass function (IMF) and evolved along prescribed evolutionary tracks (Leitherer et al. 1999). Because Sgr B1 apparently does not consist of a centralized cluster but instead contains a number of individual HII regions, each containing too few stars to completely sample the original IMF, separate modeling of each of the sub-regions can elucidate the characteristics of their ionizing stars (e.g., Simpson et al. 2004). Consequently, a set of models (plotted as triangles) was computed using the same input density, inner radii, and filling factors and using stellar SEDs computed with the WM-BASIC code (Pauldrach et al. 2001) by the Binary Population and Spectral Synthesis (BPASS) project (Eldridge et al. 2017). In Figure 3 the WM-BASIC models that have the observed values of either O⁺⁺/S⁺⁺ or [O III] 88/[S III] 33 μ m all have T_{eff} of 32,300 or 34,600 K; for the WM-BASIC code, these $T_{\rm eff}$ correspond to supergiants with spectral types O8.5-O9 and even later spectral types for dwarfs (Sternberg et al. 2003). From the lack of O stars earlier than this, we infer ages of several Myr, in agreement with the ages estimated from the comparison with Starburst99 SEDs.

On the other hand, there could possibly be contributions from high-velocity shocks. Ho et al. (2014) observed optical forbidden lines in galactic outflows in the star-forming galaxy SDSS J090005.05+000446.7. They found that their models fit the observed line ratios better if they added various amounts of shock model line intensities to the line intensities predicted by their photoionization models. Following their procedure, we plot in Figure 3(b) the effects of adding shocked-gas intensities to a sample Cloudy model. These shock models from Allen et al. (2008) are their "L_n100_b0.01" (density 100 cm^{-3} $B = 0.01 \,\mu\,\text{G}$) and "S_n1000_b0.01" (density 1000cm⁻³), $B = 0.01 \,\mu\text{G}$) model sets, plotted as cyan and brown lines, respectively. They were chosen because their neutral gas densities bracket the observed ionized gas densities and because they came closest to producing the observed [S III] $33/[Si II] 34 \,\mu m$ line ratios, compared to models with larger B. Both series of shock models have nominally solar abundance, but the line ratios were adjusted for the GC abundances from S18, which are not very different from solar. In Figure 3(b), the fractional amounts of shock model lines are added to the model overplotted with the large gray cross (a model with a WM-BASIC supergiant SED with $T_{\rm eff} =$ 30,200 K), chosen only to make it easier to distinguish the effects of adding shocks from the effects of adding X-rays

(other models could be used with similar results). Here, the connected points indicate increasing shock velocity ranging from 100, 125, ... 350 km s^{-1} (left to right), and the unconnected points (bottom to top) indicate the various fractional additions of 0.01, 0.05, 0.1, 0.3, 0.5, and 1.0 times the line intensities. As expected, the effects of adding shocks are similar to the effects of adding X-rays, no doubt because shocks produce very hot gas, which cools by radiating X-ray thermal bremsstrahlung (e.g., Allen et al. 2008).

In both panels of Figure 3, we see that the observed Ne⁺⁺/O⁺⁺ ratios or the [Ne III] 15.6/[O III] 88 μ m line ratios with few exceptions agree only with models with additional X-rays. This includes models where the X-rays originate in fast shocks (Allen et al. 2008), as seen in Figure 3(b). Shocks are especially likely for Position 1, which includes some high-energy [Ne V] emission (IP = 97 eV; Table 8 of \$18), and Positions 7 and 8, which are locations of regions with multiple radial-velocity components in the radio-recombination-line measurements of M92.

For the cluster age or stellar $T_{\rm eff}$, the plotted observations are found mostly in the vicinity of models with either Starburst99 age $=10^{6.65}-10^{6.60}$ years or $T_{\rm eff} = 32,300-34,600$ K for WM-BASIC model atmospheres (Eldridge et al. 2017). From this, we conclude that the Starburst99 age of ~4.6 Myr for Sgr B1 of S18 is not an anomaly due to the choice of line ratios used in the analysis, as we now infer a similar age with the addition of the very important O⁺⁺ ionization state.

5. Summary and Conclusions

We have mapped the GC H II region Sgr B1 with FIFI-LS on SOFIA in the lines of [O III] 52 and 88 μ m. From the extinction-corrected line ratios, we computed a generally low density, even in regions of high radio emission, with no high-density clumps indicating active or incipient massive-star formation.

We next compared the [O III] 88 μ m line intensities to intensities of [S III] and [Ne III] lines taken from archived Spitzer measurements by S18. These ionic ratios indicate that there is no central ionizing source, but instead at least eight small sub-regions with higher ionization surrounded by larger expanses of low-ionization gas. The regions of higher ionization have only a small correlation with the regions of higher density. By comparing the line ratios and ratios with H II region models computed with Cloudy (Ferland et al. 2017), we find that the results of S18 are confirmed—the region is ionized either by late O stars with $T_{\rm eff} < 35,000$ K or by stellar clusters with ages \sim 4–5 Myr (SEDs were calculated with Starburst99, Leitherer et al. 2014). X-rays in addition to those of the stellar SEDs are required to produce the measured Ne^{++}/O^{++} ratio. X-rays from fast shocks, such as those computed by Allen et al. (2008), also can produce the observed line ratios. We conclude that having lines from ions with $IP \sim 35 \, eV$ enables a significant contribution to our understanding of an H II region's ionizing SED—even the Ne⁺⁺ ionization stage (IP = 41 eV) can show the effects of local X-rays and/or shocks.

We conclude that the scattered locations of the ionizing stars and their inferred low T_{eff} , in addition to the low densities and apparent dispersal of the ionized gas (e.g., M92), all indicate that any ionizing star cluster is at least a few Myr old, in spite of the close velocity connection of the gas to the young, starforming region Sgr B2. We suggest that the stars that ionize Sgr B1 were not formed in situ, but date from a previous era of star formation and have already orbited the Galactic nucleus back to their present location at Galactic longitude $\sim 0^{\circ}.5$. This scenario is similar to the suggested origin of the Arches and Quintuplet Clusters.

Based on observations made with the NASA/DLR Stratospheric Observatory for Infrared Astronomy (SOFIA). SOFIA is jointly operated by the Universities Space Research Association, Inc. (USRA), under NASA contract NNA17BF53C, and the Deutsches SOFIA Institut (DSI) under DLR contract 50 OK 0901 to the University of Stuttgart. Financial support for this work was provided by NASA through awards 04-0113 and 05-0082 issued by USRA. We thank Christian Fischer and Randolf Klein for assistance with the observations and the referee for the thoughtful comments that improved the presentation of the Letter.

Facility: SOFIA(FIFI-LS). *Software:* Cloudy (Ferland et al. 2017).

ORCID iDs

Janet P. Simpson ^(b) https://orcid.org/0000-0001-8095-4610 Michael J. Kaufman ^(b) https://orcid.org/0000-0002-2521-1985

References

- Allen, M. G., Groves, B. A., Dopita, M. A., Sutherland, R. S., & Kewley, L. J. 2008, ApJS, 178, 20
- An, D., Ramírez, S. V., Sellgren, K., et al. 2011, ApJ, 736, 133
- An, D., Sellgren, K., Boogert, A. C. A., Ramírez, S. V., & Pyo, T. S. 2017, ApJ, 843, 36
- Barnes, A. T., Longmore, S. N., Battersby, C., et al. 2017, MNRAS, 469, 2263
- Colditz, S., Beckmann, S., Bryant, A., et al. 2018, JAI, 7, 1840004
- De Pree, C. G., Peters, T., Mac Low, M. M., et al. 2015, ApJ, 815, 123
- Eldridge, J. J., Stanway, E. R., Xiao, L., et al. 2017, PASA, 34, 58
- Ferland, G. J., Chatzikos, M., Guzmán, F., et al. 2017, RMxAA, 53, 385
- Fischer, C., Beckmann, S., Bryant, A., et al. 2018, JAI, 7, 1840003
- Habibi, M., Stolte, A., & Harfst, S. 2014, A&A, 566, A6
- Ho, I.-T., Kewley, L. J., Dopita, M. A., et al. 2014, MNRAS, 444, 3894
- Kruijssen, J. M., Dale, D., & Longmore, S. N., J. E. 2015, MNRAS, 447, 1059
- Lang, C. C., Goss, W. M., Cyganowski, C., & Clubb, K. I. 2010, ApJS, 191, 275
- Lang, C. C., Goss, W. M., & Morris, M. 2002, AJ, 124, 2677
- Leitherer, C., Ekström, S., Meynet, G., et al. 2014, ApJS, 212, 14
- Leitherer, C., Schaerer, D., Goldader, J. D., et al. 1999, ApJS, 123, 3
- Longmore, S. N., Kruijssen, J. M. D., Bally, J., et al. 2013, MNRAS, 433, L15
- Mauerhan, J. C., Muno, M. P., Morris, M. R., Stolovy, S. R., & Cotera, A. 2010, ApJ, 710, 706
- Mehringer, D. M., Palmer, P., & Goss, W. M. 1993, ApJL, 402, L69
- Mehringer, D. M., Palmer, P., & Goss, W. M. 1995, ApJS, 97, 497
- Mehringer, D. M., Yusef-Zadeh, F., Palmer, P., & Goss, W. M. 1992, ApJ, 401, 168 (M92)
- Molinari, S., Bally, J., Noriega-Crespo, A., et al. 2011, ApJL, 735, L33
- Molinari, S., Scisano, E., Elia, D., et al. 2016, A&A, 591, A149
- Pauldrach, A. W. A., Hoffmann, T. L., & Lennon, M. 2001, A&A, 375, 161
- Rubin, R. H., Simpson, J. P., Colgan, S. W. J., et al. 2008, MNRAS, 387, 45
- Schneider, F. R. N., Izzard, R. G., de Mink, S. E., et al. 2014, ApJ, 780, 117 Simpson, J. P. 2018, ApJ, 857, 59 (S18)
- Simpson, J. P., Colgan, S. W. J., Cotera, A. S., et al. 2007, ApJ, 670, 1115
- Simpson, J. P., Cotera, A. S., Burton, M. G., et al. 2012, MNRAS, 419, 211
- Simpson, J. P., Rubin, R. H., Colgan, S. W. J., Erickson, E. F., & Haas, M. R. 2004, ApJ, 611, 338
- Sternberg, A., Hoffmann, T. L., & Pauldrach, A. W. A. 2003, ApJ, 599, 1333 Stolte, A., Hußmann, B., Morris, M. R., et al. 2014, ApJ, 789, 115
- Storey, P. J., & Sochi, T. 2015, MNRAS, 449, 2974
- Temi, P., Marcum, P. M., Young, E., et al. 2014, ApJS, 212, 24
- Young, E. T., Becklin, E. E., Marcum, P. M., et al. 2012, ApJL, 749, L17